

## RHEOLOGICAL PROPERTIES OF LATE-STAGE LAVA FLOWS ON ASCRAEUS MONS: NEW EVIDENCE FROM HRSC. H. Hiesinger<sup>1,2</sup>, J.W. Head III<sup>1</sup>, G. Neukum<sup>3</sup> and the HRSC Co-Investigator Team.

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**Introduction:** The Tharsis Montes, Arsia Mons, Pavonis Mons, and Ascraeus Mons, are large volcanic constructs that are part of the Tharsis rise. The Tharsis rise is commonly interpreted to be the result of a long-lasting large mantle upwelling that constructed the rise through a combination of uplift and magmatism (plutonism and volcanism) [e.g., 1-7]. Of particular interest is the construction of the huge individual edifices, their ages, duration, episodicity [8], and rheology. Here we report on estimates of the rheological properties of late-stage lava flows on the eastern flank of Ascraeus Mons, Mars.

**Data:** Our estimates are based on new images obtained by the High Resolution Stereo Camera (HRSC) on board ESA's Mars Express spacecraft. The HRSC camera is a linescan camera with 9 CCD lines (blue green, red, IR, 3 stereo channels, 2 photometric channels) oriented perpendicular to the flight direction. The HRSC camera acquires images at spatial resolutions of about 10 m/pixel and is complemented by a Super Resolution Channel (SRC) with a 1024 x 1032 pixel frame CCD, which obtains images of about 2.3 m/pixel from an altitude of 250 km at periaapsis. We used HRSC data from orbit 16 to investigate lava flows on Ascraeus Mons, one of the three Tharsis Montes. Compared to earlier studies, the high spatial resolution of the HRSC data allowed us to map 25 late-stage lava flows and to measure their dimensions, as well as their morphological characteristics in greater detail. Values for the slope  $\alpha$ , flow length  $l$ , flow width  $w$ , flow height  $h$ , total levee width  $w_l$ , and the width of the leveed flow channel  $w_c$  can be determined directly either from MOLA data or from the imaging data. Length and width of these lava flows can be readily measured on HRSC images, but the height of these flows is currently below the vertical resolution of digital elevation models (DEM) that can be confidently calculated from the HRSC data at this time. Such height measurements will be possible in the near future when more experience has been gained in producing high-resolution DEMs. Here we used shadow measurements along the lava flows and the flow fronts.

**Morphology and Dimensions:** In the HRSC images we observe several lava flows with well-defined leveed channels, some of which are truncated by the collapse of the calderas and extend for tens of kilometers downslope. We measured the length of

each flow and performed about 184 measurements of the flow width, 224 measurements of the flow height, and 80 measurements for the levee and channel widths. We find the average flow length to be  $\sim 19$  km, ranging from  $\sim 4.1$  to  $\sim 38.3$  km. The average width is  $\sim 1.3$  km, with a minimum of  $\sim 570$  m and a maximum of  $\sim 2030$  m. We find the average leveed flow width to be on the order of  $\sim 990$  m, ranging from  $\sim 590$  m to  $\sim 2270$  m. The average levee width varies between  $\sim 450$  m and  $\sim 1740$  m. with an average of  $\sim 710$  m. The average channel width is about 280 m, with a minimum of  $\sim 140$  m and a maximum of  $\sim 530$  m. We performed 224 individual shadow measurements along the 25 investigated flows. We find that these flows are on average 39 m thick, varying from  $\sim 25$  to  $\sim 90$  m. These results are in good agreement with measurements of 25-220 m [9], 5-65 m [10],  $\sim 30$  m [11],  $65 \pm 20$  m [12], 13-33 m [13], and  $\sim 20$  m [14] and greater than reported from Olympus Mons (4-26m) [28]. The average slopes of the studied flows range from  $\sim 1.5^\circ$  to  $\sim 6.7^\circ$ , with most flows being emplaced on slopes of  $\sim 3.5^\circ$ .

On the basis of morphologic similarities between terrains on Ascraeus Mons and terrestrial shield volcanoes, *Zimbelman and McAllister* [15] proposed that individual prominent flows on Ascraeus Mons are a'a flows and the planar areas adjacent to the flows are pahoehoe flows. *Wilson et al.* [16] reported that if no other factors intervene, thermal constraints will be the limiting factor for the maximum length of a flow fed by a given volume or mass effusion rate. They classified lava flows into several categories, including (1) cooling-limited flows, (2) volume-limited flows, (3) accidentally-breached flows, (4) break-out flows, (5) captured flows, (6) tube-fed flows. We find that our flow N shares numerous characteristics of a volume-limited flow such as short length, a drained channel, and no break-out flows. However the majority of flows appear to be more akin to cooling-limited flows or break-out flows. Flows E1 - E3 might have formed as accidentally-breached flows or break-out flows. We did not find evidence for captured flows and evidence for tube-fed flows remains at best ambiguous.

**Yield strength:** Lava flows are often modeled as a Bingham plastic controlled by two parameters, the yield strength and the plastic viscosity [e.g., 17]. The

yield strength  $\tau$  of lava flows can be related to the flow dimensions by following equations

$$\begin{aligned}\tau &= \rho g \sin\alpha h \\ \tau &= \rho g h^2/w \\ \tau &= \rho g \sin^2\alpha 2w_1 \\ \tau &= \rho g \sin^2\alpha (w-w_c)\end{aligned}$$

Gravity  $g$  is known as  $3.7278 \text{ ms}^{-2}$  and density  $\rho$  was chosen to be  $2,500 \text{ kgm}^{-3}$ . Other input parameters, e.g., the flow height and width, channel and levee width, and the slope angle, were derived from HRSC or MOLA data. Depending on the equation used, we find a minimum yield strength of  $\sim 5.1 \times 10^3$  and a maximum yield strength of  $\sim 1.2 \times 10^5 \text{ Pa}$ . Calculating the average of all derived yield strengths, we find a value of  $\sim 2.7 \times 10^4 \text{ Pa}$ ; a yield strength very similar to that published by *Zimbelman* [11], i.e.,  $\sim 2.1 \times 10^4 \text{ Pa}$ . These values are comparable to estimates for terrestrial basaltic lava flows, and are in good agreement with estimates of *Zimbelman* [11] derived for a small number of lava flows on Ascræus Mons.

**Effusion rate:** The effusion rate  $Q$  is given as

$$Q = G_z \kappa x w/h$$

where  $G_z$  is the dimensionless Graetz number,  $\kappa$  is the thermal diffusivity ( $\text{m}^2\text{s}^{-1}$ ),  $x$  is the flow length (m), and  $w$  and  $h$  is defined as above. In analogy to terrestrial lava flows we assumed a Graetz number of 300 and a thermal diffusivity of  $3 \times 10^{-7} \text{ m}^2\text{s}^{-1}$ . As a result we find that effusion rates range from  $\sim 7$  to  $\sim 210 \text{ m}^3\text{s}^{-1}$ , averaging about  $70 \text{ m}^3\text{s}^{-1}$ . Such effusion rates for the Ascræus flows studied are similar to effusion rates of  $18\text{--}60 \text{ m}^3\text{s}^{-1}$ , with an average of  $35 \text{ m}^3\text{s}^{-1}$  [11].

**Emplacement time:** On the basis of our estimates of the effusion rates and the measured dimensions of the flows, we calculated that the time necessary to emplace the flows is on average on the order of tens to hundreds of days.

**Viscosity:** For the determination of the viscosities  $\eta$  we made use of the following equation [e.g., 12, 18]

$$h = (Q \eta / \rho g)^{1/4}$$

This equation assumes a Newtonian flow behavior and is therefore in conflict with the assumption that lava flows have a Bingham rheology [e.g., 19–20].

Jeffrey's equation also relates the viscosity of a flow to its effusion rate and its dimensions [e.g., 21–22].

$$\eta = (\rho g h^3 w \sin\alpha) / nQ$$

In this equation  $n$  is a constant equal to 3 for broad flows and 4 for narrow flows. *Gregg and Fink* [22] pointed out that although Jeffrey's equation has been widely used to derive lava flow characteristics, it requires the questionable assumption that lava behaves as a Newtonian fluid.

*Wilson and Head* [17] and *Zimbelman* [11] presented another method to calculate viscosities

$$\begin{aligned}r &= w_c / (w - w_c) \\ \eta &= (w_c^3 \tau \sin^2\alpha) / (24Q) \quad \text{for } r < 1 \\ \eta &= (w_c^{11/4} \tau^{5/4} \sin^{6/4}\alpha) / (24Q g^{1/4} \rho^{1/4}) \quad \text{for } r \geq 1\end{aligned}$$

As a result we calculated average viscosities of  $\sim 5 \times 10^5$  to  $\sim 6 \times 10^6 \text{ Pa}\cdot\text{s}$ . Minimum viscosities are on the order of  $3 \times 10^4 \text{ Pa}\cdot\text{s}$ , maximum viscosities are about  $4 \times 10^7 \text{ Pa}\cdot\text{s}$ . Using Jeffrey's equation yields significantly higher average viscosities of  $6.2 \times 10^8 \text{ Pa}\cdot\text{s}$ , ranging from  $1.7 \times 10^7$  to  $4.2 \times 10^9 \text{ Pa}\cdot\text{s}$ . The values based on Jeffrey's equation appear to be unrealistically high compared to basalt/andesite flows on Earth and other planetary bodies. Our results are in excellent agreement with previously published viscosities, which range from  $9.7 \times 10^5$  to  $6.9 \times 10^6 \text{ Pa}\cdot\text{s}$  [e.g., 11, 12, 20, 23]. In addition, our results are consistent with viscosities of terrestrial basalts and andesites, which are on the order of  $1.4 \times 10^2$  to  $1 \times 10^7 \text{ Pa}\cdot\text{s}$  [e.g., 20, 24–27].

For the future we plan to include older lava flows on Ascræus Mons in order to investigate possible changes in the rheologic properties with time [e.g., 11]

**Conclusions:** In summary, the new data allow us to estimate more precisely the areas and volumes of lava flows, the duration and rates of their emplacement, their petrologic evolution, their relation to the geologic history of Mars and their influence on the atmosphere and the environment. On the basis of our investigation we conclude that the lava flows investigated are likely to be basaltic to andesitic in composition. A comparison to terrestrial flows indicates that the investigated flows share several characteristics with terrestrial a'a flows.

**References:** [1] Banerdt et al. (1992) in *Mars*, Univ. of Arizona Press; [2] Breuer et al. (1996) *J. Geophys. Res.* 101; [3] Harder (1998) *J. Geophys. Res.* 103; [4] Smith et al. (1999) *Science* 284; [5] Zuber et al. (2000) *Science* 28; [6] Head and Solomon (1981) *Science* 213; [7] Phillips et al. (2001) *Science* 291; [8] Neukum et al. (2004) *Nature* 432; [9] Head et al. (1998) *LPSC XXIX*, #1125; [10] Schaber et al. (1978) *PLPSC* 9; [11] Zimbelman (1985) *PLPSC* 16; [12] Warner and Gregg (2003) *J. Geophys. Res.*, 108; [13] Peitersen et al. (2001) *LPSC XXXII*, #1472; [14] Glaze et al. (2003) *LPSC XXXIV*, #1315; [15] Zimbelman and McAllister (1985) *LPSC XVI*; [16] Wilson et al. (1993) *LPSC XXIV*; [17] Wilson and Head (1983) *Nature* 302; [18] Fink and Griffiths (1990) *J. Fluid Mech.*, 22; [19] Shaw et al. (1968) *Amer. J. Sci.*, 266; [20] Hulme (1976) *Icarus*, 27; [21] Gregg and Zimbelman (2000) in *Environmental Effects on Volcanic Eruptions: From Deep Oceans to Deep Space*; [22] Gregg and Fink (1996) *J. Geophys. Res.*, 101; [23] Cattermole (1987) *PLPSC* 17; [24] Moore (1987) *U.S. Geol. Surv. Prof. Paper* 1350; [25] Cigolini et al. (1984) *J. Volcanol. Geotherm. Res.*, 29; [26] Pinkerton and Sparks (1976) *J. Volcanol. Geotherm. Res.*, 1; [27] Murase and McBirney (1973) *Geol. Soc. Am. Bull.*, 84. Basilevskaya (2005) *LPSC XXXVI*, #1082.